

Fluctuating Hydrodynamics of Flow through Porous Membranes

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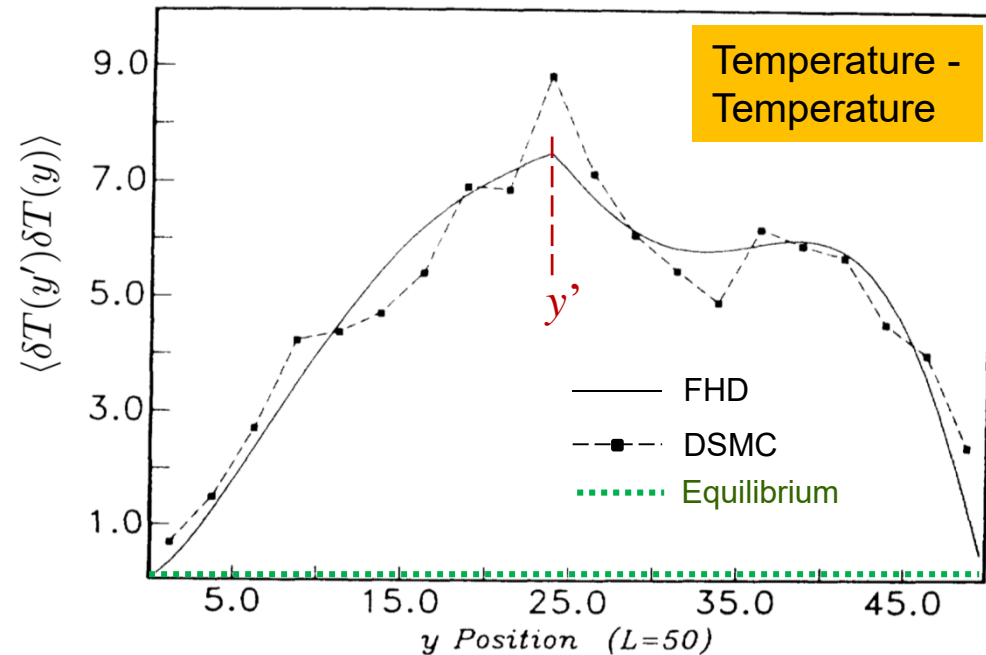
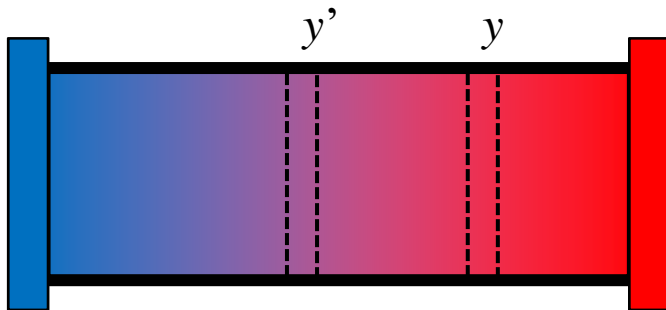


Non-equilibrium Fluctuations

Hydrodynamic fluctuations are long-ranged in a fluid held at a non-equilibrium steady state.

$$\langle \delta\rho(y')\delta v(y) \rangle \propto \nabla T$$

$$\langle \delta T(y')\delta T(y) \rangle \propto (\nabla T)^2$$



M. Malek Mansour, ALG, G. Lie and E. Clementi, *Phys. Rev. Lett.* **58** 874 (1987).

Fluctuating Hydrodynamics (FHD)

Landau and Lifshitz introduced stochastic flux terms into the equations of hydrodynamics to model spontaneous fluctuations in fluids.

$$\begin{aligned}\frac{\partial}{\partial t}(\rho) &= -\nabla \cdot (\rho \mathbf{u}), \\ \frac{\partial}{\partial t}(\rho \mathbf{u}) &= -\nabla \cdot (\rho \mathbf{u} \otimes \mathbf{u}) - \nabla P - \nabla \cdot [\boldsymbol{\Pi} + \tilde{\boldsymbol{\Pi}}], \\ \frac{\partial}{\partial t}(\rho E) &= -\nabla \cdot (\rho \mathbf{u} E + P \mathbf{u}) - \nabla \cdot [\mathbf{Q} + \tilde{\mathbf{Q}}] - \nabla \cdot ([\boldsymbol{\Pi} + \tilde{\boldsymbol{\Pi}}] \cdot \mathbf{u}).\end{aligned}$$

[Stress tensor]

[Heat flux]

Dissipative Fluxes in FHD

The deterministic stress tensor and heat flux take their standard linear forms (Stokes and Fourier laws),

$$\Pi_{ij} = -\eta \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right) + \delta_{ij} \left(\frac{2}{3} \eta \nabla \cdot \mathbf{u} \right), \quad \text{and} \quad \mathbf{Q} = -\kappa \nabla T$$

Stochastic stress tensor and heat flux are independent noises, white in space and time, with zero mean and variances,

$$\langle \tilde{\Pi}_{ij}(\mathbf{r}, t) \tilde{\Pi}_{kl}(\mathbf{r}', t') \rangle = 2k_B \eta T \left[(\delta_{ik} \delta_{jl} + \delta_{il} \delta_{jk}) - \frac{1}{3} \delta_{ij} \delta_{kl} \right] \delta(t - t') \delta(\mathbf{r} - \mathbf{r}'),$$

$$\langle \tilde{Q}_i(\mathbf{r}, t) \tilde{Q}_j(\mathbf{r}', t') \rangle = 2k_B \kappa T^2 \delta_{ij} \delta(t - t') \delta(\mathbf{r} - \mathbf{r}').$$

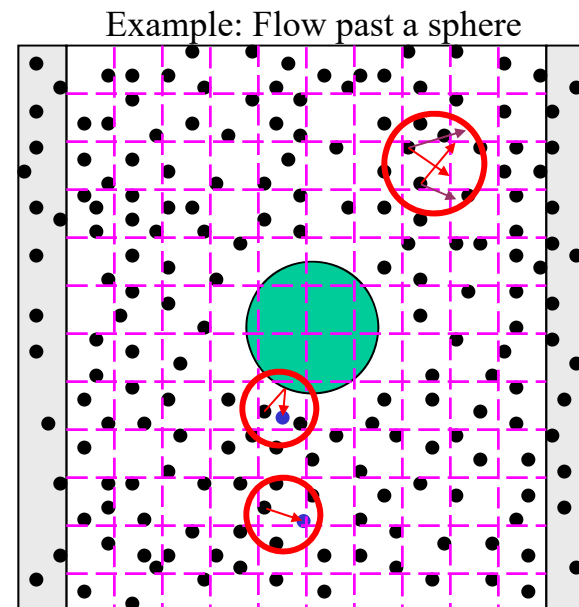
Finite Volume SPDE solver

Write the FHD equations as: $\frac{\partial}{\partial t} \mathbf{U} = -\nabla \cdot \mathbf{F} - \nabla \cdot \mathbf{D} - \nabla \cdot \tilde{\mathbf{S}}$

- Time integration: $\frac{\partial}{\partial t} \mathbf{U} = \frac{\partial}{\partial t} (\rho, \rho \mathbf{u}, \rho E)$ • Three-stage Runge-Kutta
- Hyperbolic: $\mathbf{F} = (\rho \mathbf{u}, \rho \mathbf{u} \mathbf{u} + P \mathbf{I}, \rho \mathbf{u} E + \mathbf{u} P)$ • Four point centered
- Parabolic: $\mathbf{D} = (0, \mathbf{\Pi}, \mathbf{Q} + \mathbf{u} \cdot \mathbf{\Pi})$ • Two point centered
- Stochastic: $\tilde{\mathbf{S}} = (0, \tilde{\mathbf{\Pi}}, \tilde{\mathbf{Q}} + \mathbf{u} \cdot \tilde{\mathbf{\Pi}})$ • Weighted 2 point centered

Direct Simulation Monte Carlo (DSMC)

- Initialize system with particles
- Loop over time steps
 - Create particles at open boundaries
 - Move *all* the particles
 - Process any interactions of particle & boundaries
 - Sort particles into cells
 - Sample statistical values
 - Select and execute random collisions

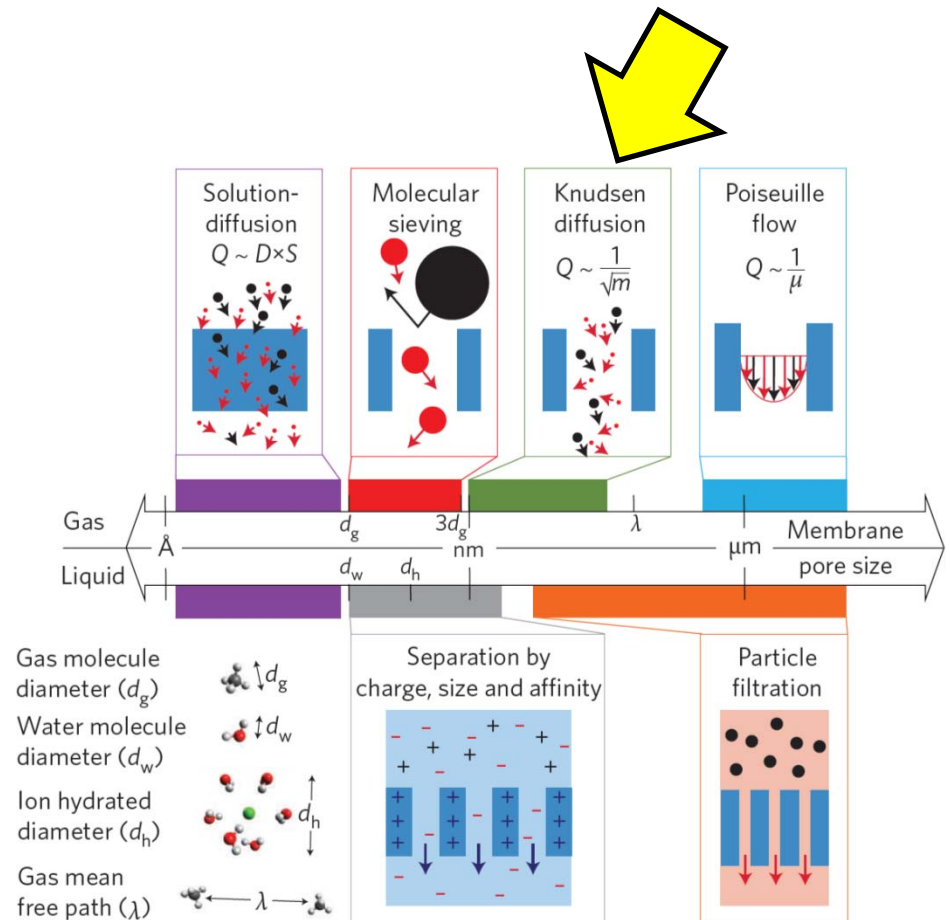


G.A. Bird, *Molecular Gas Dynamics and Direct Simulation of Gas Flows*, Clarendon, Oxford (1994)
F. Alexander and ALG, *Computers in Physics*, **11** 588 (1997)

Porous Membranes

Effusion (Knudsen diffusion) is gas transport through a membrane with pore sizes roughly between the mean free path and the molecule diameter.

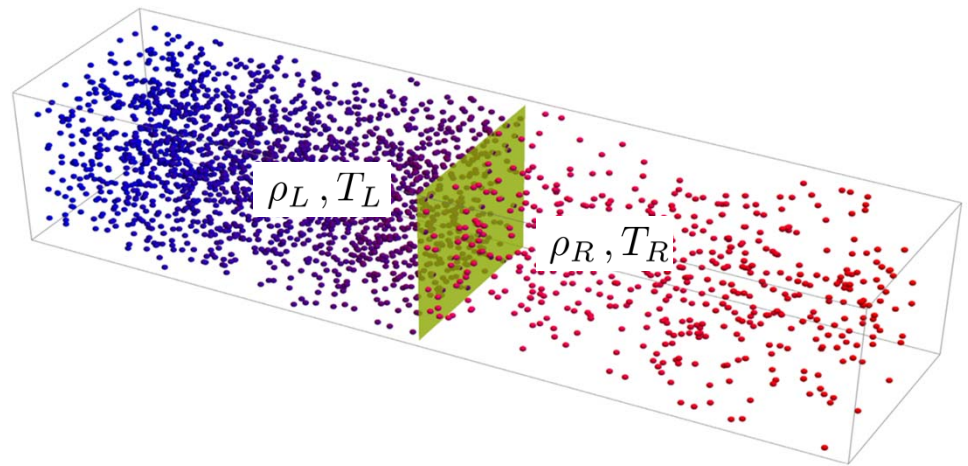
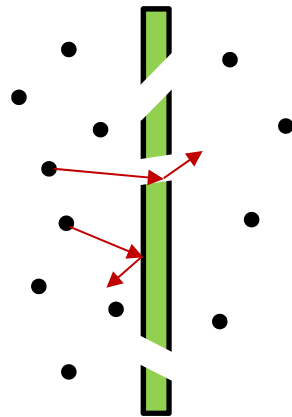
For standard conditions this range in pore size is between 100 nm to 2 nm.



L. Wang, et al., *Nature Nanotechnology* **12** 509–522 (2017)

Transport by Effusion

Molecules reaching the interface cross it with probability f .



The mean fluxes of mass and energy are,

$$\langle J_M \rangle = \frac{f A k_B^{1/2}}{\sqrt{2\pi m}} \left(\rho_L T_L^{1/2} - \rho_R T_R^{1/2} \right), \quad \langle J_E \rangle = \frac{2f A k_B^{3/2}}{m \sqrt{2\pi m}} \left(\rho_L T_L^{3/2} - \rho_R T_R^{3/2} \right).$$

Langevin Model for Effusion Membrane

In FHD we model the mass and energy crossing the membrane with the Langevin equations,

$$\frac{d}{dt}M = \langle J_M \rangle + \tilde{J}_M, \quad \frac{d}{dt}\mathcal{E} = \langle J_{\mathcal{E}} \rangle + \tilde{J}_{\mathcal{E}},$$

where the white noises have variances and covariances,

$$\langle \tilde{J}_M(t) \tilde{J}_M(t') \rangle = \frac{mfA k_B^{1/2}}{\sqrt{2\pi m}} \left(\rho_R T_R^{1/2} + \rho_L T_L^{1/2} \right) \delta(t - t')$$

$$\langle \tilde{J}_{\mathcal{E}}(t) \tilde{J}_{\mathcal{E}}(t') \rangle = \frac{6fA k_B^{5/2}}{m\sqrt{2\pi m}} \left(\rho_R T_R^{5/2} + \rho_L T_L^{5/2} \right) \delta(t - t'),$$

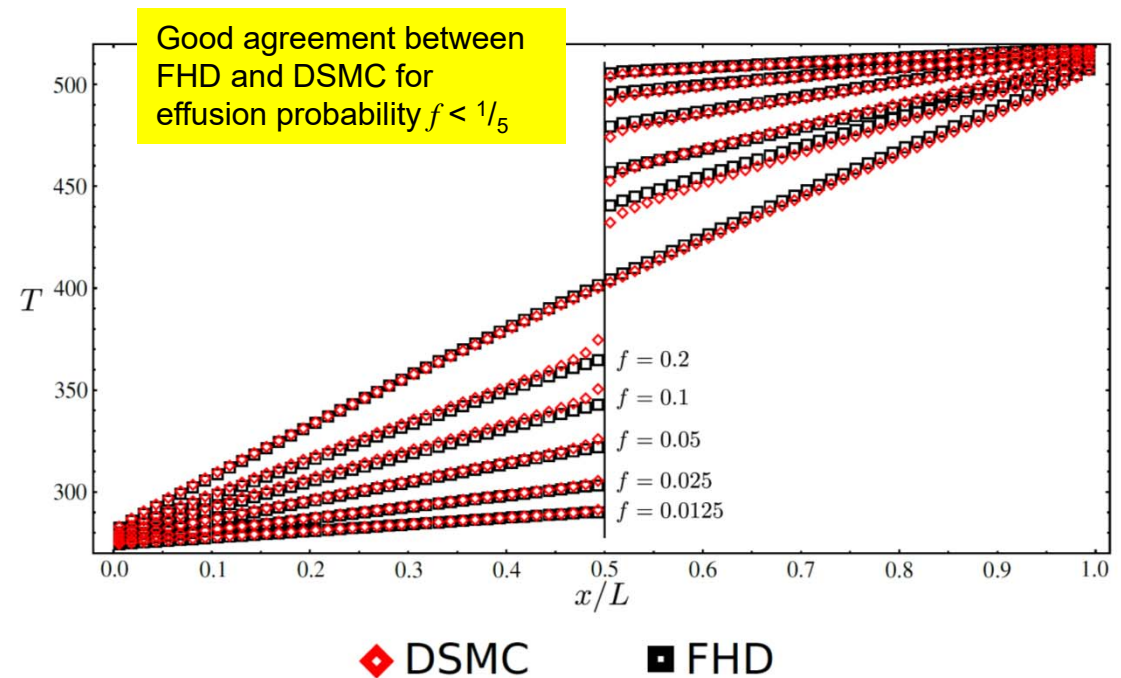
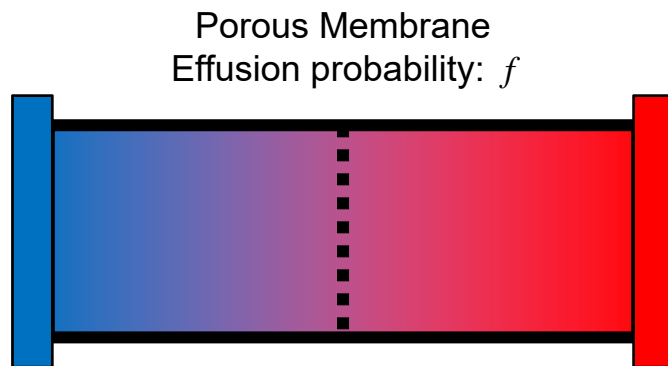
$$\langle \tilde{J}_{\mathcal{E}}(t) \tilde{J}_M(t') \rangle = \frac{2fA k_B^{3/2}}{\sqrt{2\pi m}} \left(\rho_R T_R^{3/2} + \rho_L T_L^{3/2} \right) \delta(t - t').$$



Note: Mass and energy noises are correlated

Temperature Profiles

Simulated a dilute gas with a temperature gradient in a system bisected by a porous effusion membrane.

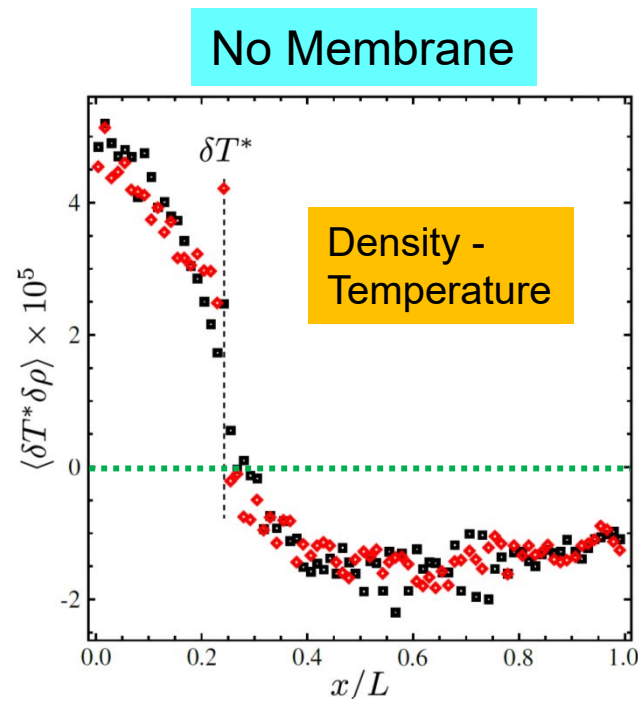


D. Ladiges, A. Nonaka, J.B. Bell, and ALG,
Physics of Fluids **31**, 052002 (2019)

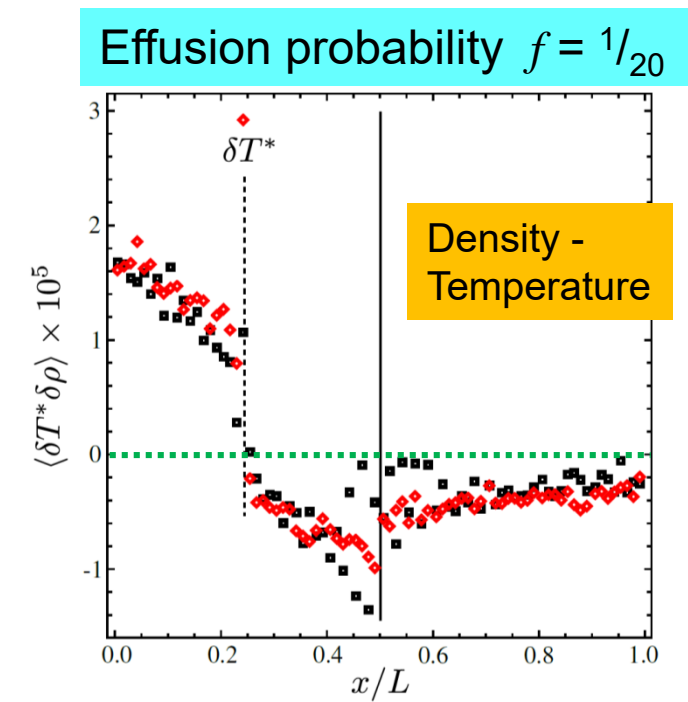
Correlations of Fluctuations

Temperature gradient is reduced by ~ 3 and so is density-temperature correlation

Correlation persists across the membrane



◆ DSMC



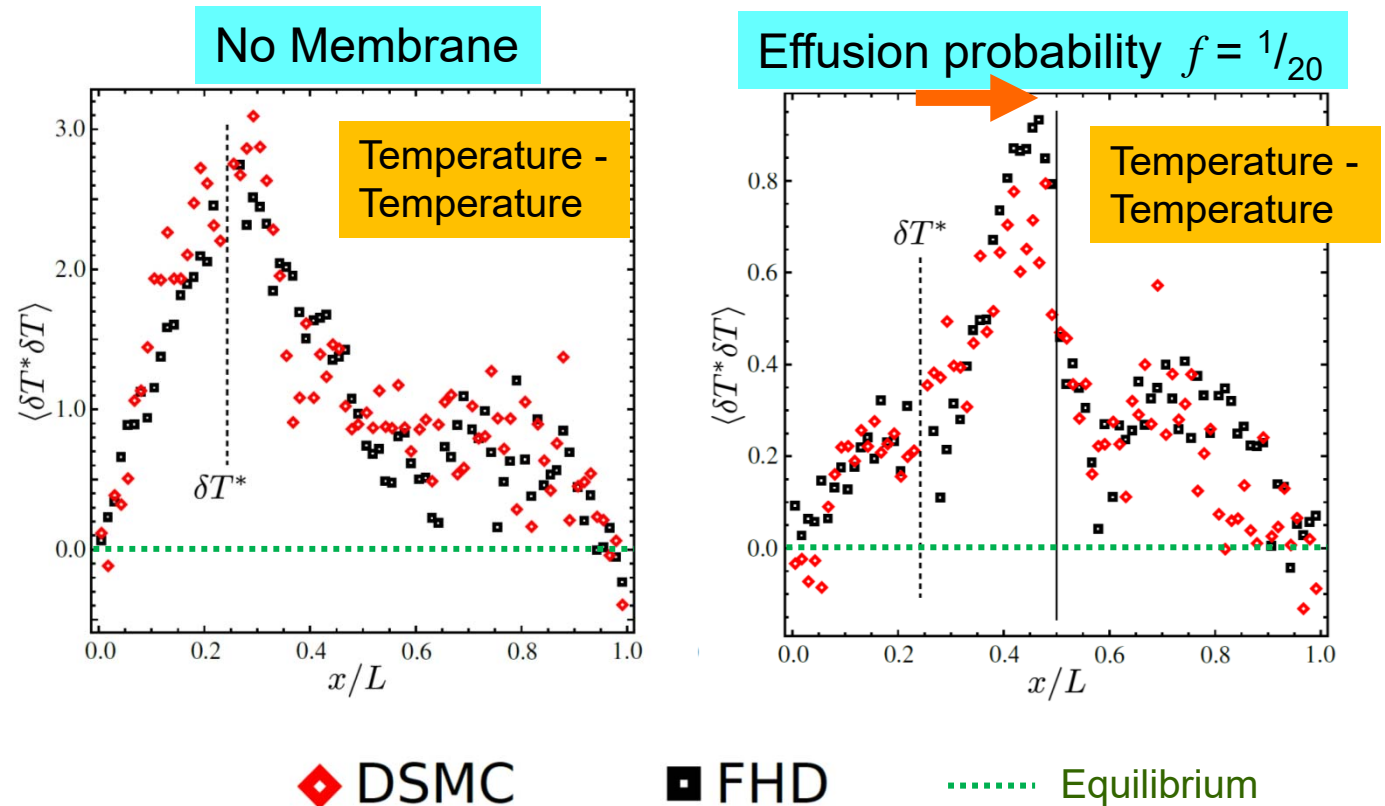
■ FHD

..... Equilibrium

Correlations of Fluctuations

Temperature-temperature correlation is reduced but persists across the membrane

Correlation peak is significantly shifted towards the membrane



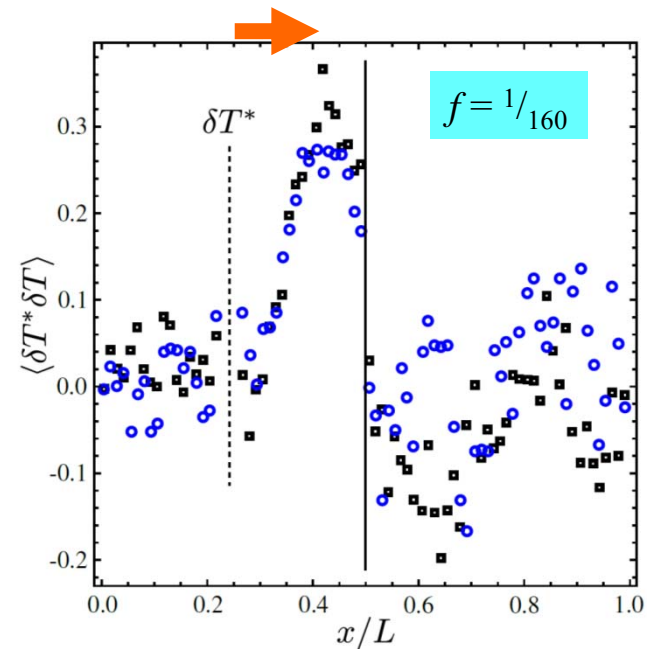
Master Equation Interface Model

Also tested a Master equation formulation using the Gillespie (SSA) algorithm in FHD.

$$\tau_e = \frac{m}{\rho f A} \sqrt{\frac{2\pi m}{k_B T}} \quad \text{Mean waiting time for crossings}$$

$$P_\epsilon(\epsilon) = \frac{\epsilon}{k_B^2 T^2} e^{-\epsilon/(k_B T)} \quad \text{Distribution of molecule energies}$$

The Langevin and Master equation models produced equivalent results.



○ Langevin equation ■ Master equation

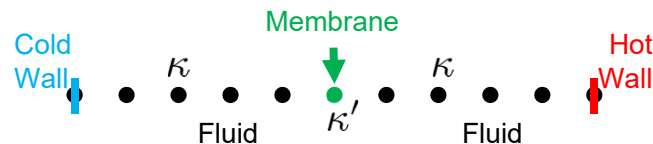
Stochastic Heat Equation

To investigate the shift in the peak we consider the linearized stochastic heat equation,

$$\rho c_v \frac{\partial}{\partial t} \delta T(y, t) = \frac{\partial}{\partial y} \kappa(y) \frac{\partial}{\partial y} \delta T - \frac{\partial}{\partial y} \tilde{Q} \quad \text{with} \quad \langle \tilde{Q}(y, t) \tilde{Q}(y', t') \rangle = 2k_B \kappa(y) T_0(y)^2 \delta(t - t') \delta(y - y').$$

Discretizing in space, $\mathbf{U} = [\delta T_1, \dots, \delta T_N]$

$$\frac{\partial}{\partial t} \mathbf{U} = -\mathbf{A} \mathbf{U} + \mathbf{B} \boldsymbol{\xi}$$



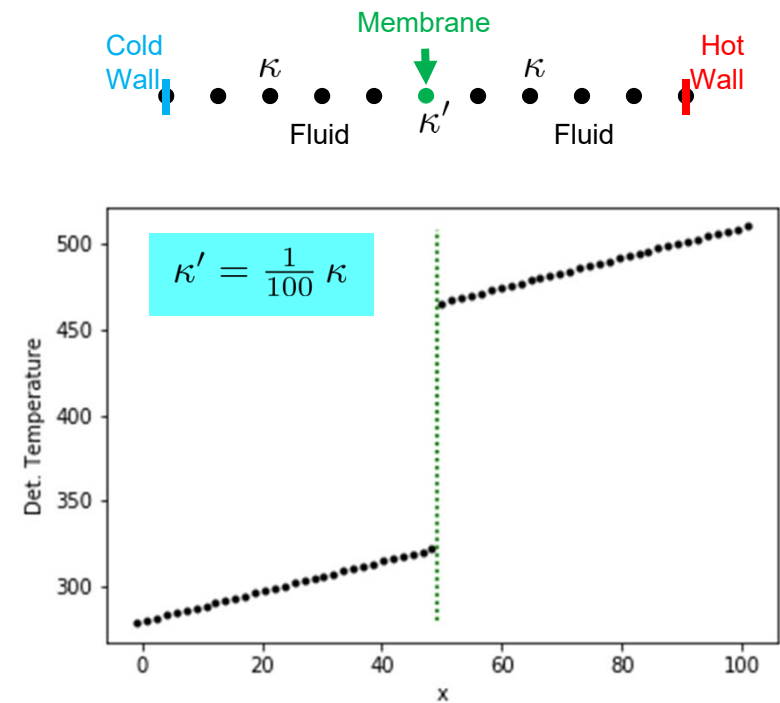
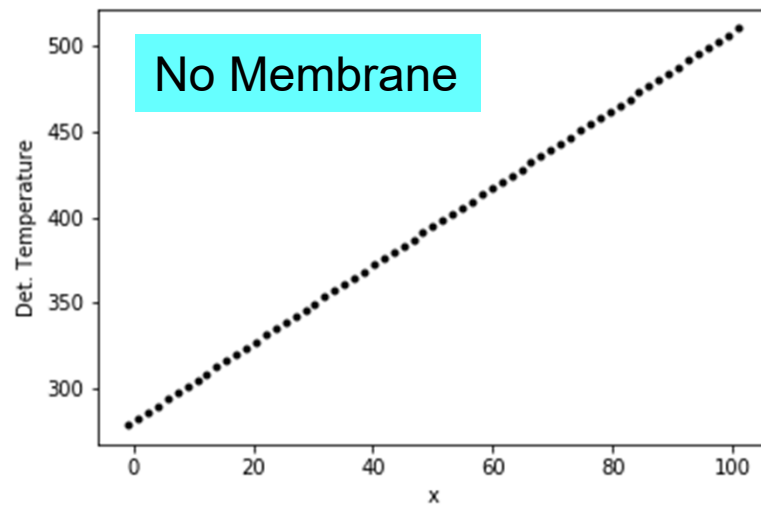
For this Ornstein-Uhlenbeck process we find the covariance $\sigma_{i,j} = \langle \delta T_i \delta T_j \rangle$ by solving

$$\mathbf{A} \boldsymbol{\sigma} + \boldsymbol{\sigma} \mathbf{A}^T = \mathbf{B} \mathbf{B}^T$$

using linear algebra (e.g., by numerical relaxation)

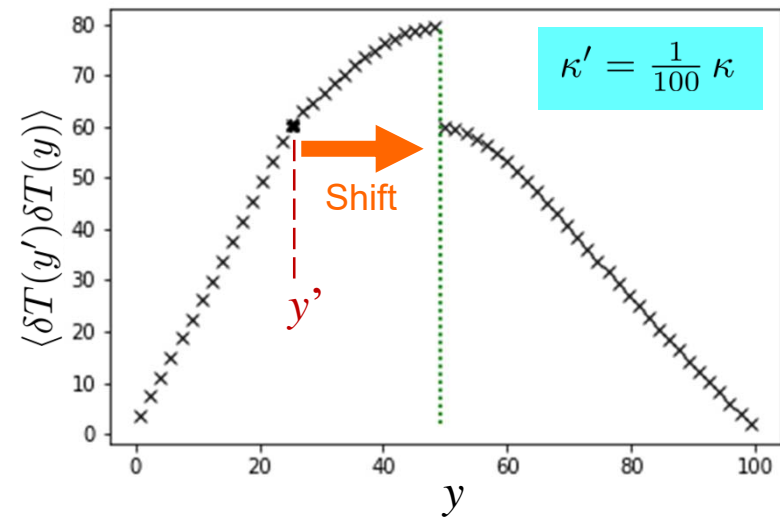
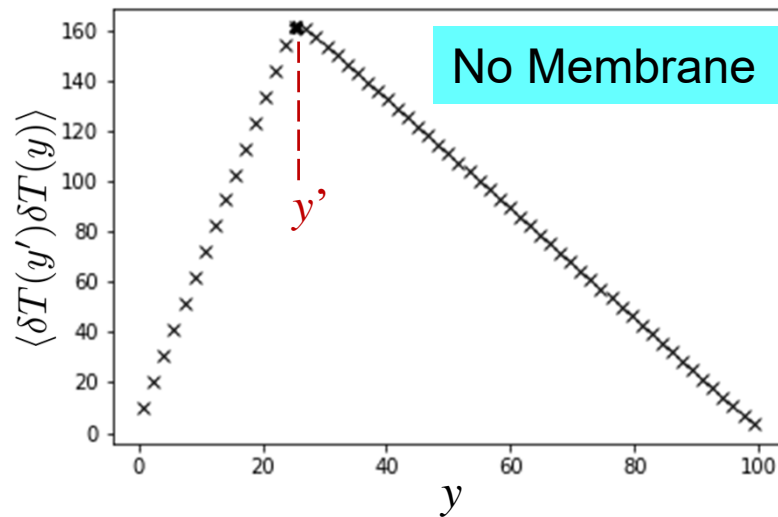
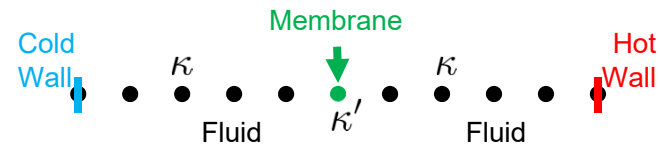
Temperature Profile

Deterministic temperature profiles are qualitatively similar to FHD and DSMC results



Correlations of Fluctuations

Temperature-temperature correlation is qualitatively similar to the FHD and DSMC result.



Summary & Future Work

Summary:

- Long range correlations persist through an effusive interface.
- Reduced magnitude largely due to change in ∇T .
- Distortion predicted by the stochastic heat equation.
- FHD can simulate gas transpiration; faster than DSMC.

Future work:

- Shear gradient; Concentration gradient
- Correlations parallel to the membrane
- Molecular sieving
- Ion transport in electrolytes
- Active transport (transport & chemistry)

Thank you for your attention

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Thank you Aleks Donev for organizing this session